

High-Quality Dual-Plasmonic Au@Cu_{2-x}Se Nanocrescents with Precise Cu_{2-x}Se Domain Size Control and Tunable Optical Properties in the Second Near-Infrared Biowindow

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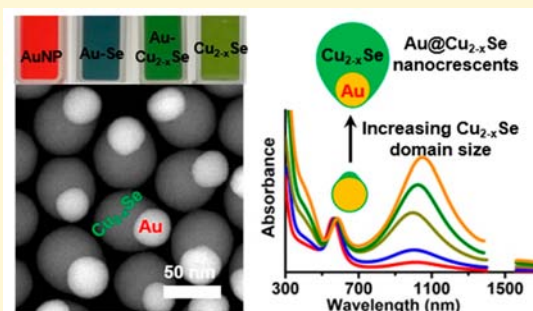
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Supporting Information

ABSTRACT: Dual-plasmonic Au–Cu_{2-x}E ($x > 0$, E = S, Se, Te) hybrids present multifunctionalities for broad applications in photovoltaics, photocatalysis, and biomedicine. However, despite the increasing interest in these dual-plasmonic Au–Cu_{2-x}E hybrids, precise control over the morphology and size of their substituents is crucial but greatly challenging. Herein, a selenium-mediated two-step synthetic approach is developed to prepare the Au@Cu_{2-x}Se nanocrescents with precise control over their size and tunable plasmonic properties through varying the concentration of the SeO₂ precursor. The complete encapsulation of the Au domain within the Cu_{2-x}Se layer in the nanocrescent in an asymmetric manner enables the strong coupling between their substituents, significantly enhancing their optical properties supported by both experimental results and theoretical modeling. The increasing Cu_{2-x}Se domain size enables the tuning of carrier densities and thereby molar extinction coefficients corresponding to the respective visible and near-infrared (NIR) plasmon bands of the Au@Cu_{2-x}Se nanocrescents. Electron transfer from the Au to Cu_{2-x}Se domains is unambiguously determined in the Au@Cu_{2-x}Se nanocrescents by transient absorption (TA) and extinction spectral measurements. The Cu_{2-x}Se-dependent bleaching amplitude and carrier dynamics of both visible and NIR plasmon TA signals are observed, along with variations of the carrier–phonon and phonon–phonon relaxation processes. This work not only presents a facile synthetic approach of high-quality dual-plasmonic Au@Cu_{2-x}Se nanocrescents and a fundamental understanding of the coupling mechanism in the Au–Cu_{2-x}Se hybrid system, but also provides a paradigm for the study on the interplay in a dual-plasmonic metal–semiconductor system.



INTRODUCTION

Plasmonic materials find a plethora of applications ranging from photovoltaics and light-driven catalysis to biomedicine such as biosensing, bioimaging, and cancer theranostics.^{1–4} Noble metals and nonstoichiometric copper chalcogenides (Cu_{2-x}E, $x > 0$, E = S, Se, Te) are among the most studied plasmonic materials and have been extensively exploited as building blocks of multifunctional materials for those plasmon-based applications.^{5–9} Au and Ag are the most commonly used plasmonic noble metals, whose surface plasmon resonances (SPRs) mainly arise from the collective oscillation of their free conduction electrons driven by the electromagnetic field of incident light.^{5,10} The large intrinsic free carrier (electron) density ($\sim 10^{23}$ cm⁻³) in noble metals makes their SPR bands in the visible (400–700 nm) and first near-infrared (NIR-I, 700–900 nm) regions;¹¹ the plasmonic response of these noble metals strongly depends on their composition, particle size, and geometric shape, and the dielectric function of the surrounding environments.^{2,12–14} Although tailoring the

anisotropic shape (i.e., nanorods, nanocages, hollow spheres, and core–shells) allows for tuning the SPR band of noble metal nanostructures into the second NIR biowindow (NIR-II, 1000–1350 nm) favorable for in vivo biomedical applications, both large particle size of >100 nm and high power density exceeding the maximum laser permissible exposure (0.33 W/cm² for 808 nm laser, 1.0 W/cm² for 1064 nm laser) restrict their in vivo bioapplications.^{15–17} Nonstoichiometric copper chalcogenides are a novel class of p-type semiconducting materials with characteristic SPR bands in the NIR spectral region especially in the NIR-II transparent biowindow. Unlike noble metals, their plasmons are supported by the positively charged holes as a result of the Cu vacancy, with a carrier density of 10¹⁹–10²¹ cm⁻³ relatively lower than noble metals.^{18,19} The SPR band of Cu_{2-x}E can be effectively

Received: October 7, 2019

Revised: November 13, 2019

Published: November 13, 2019

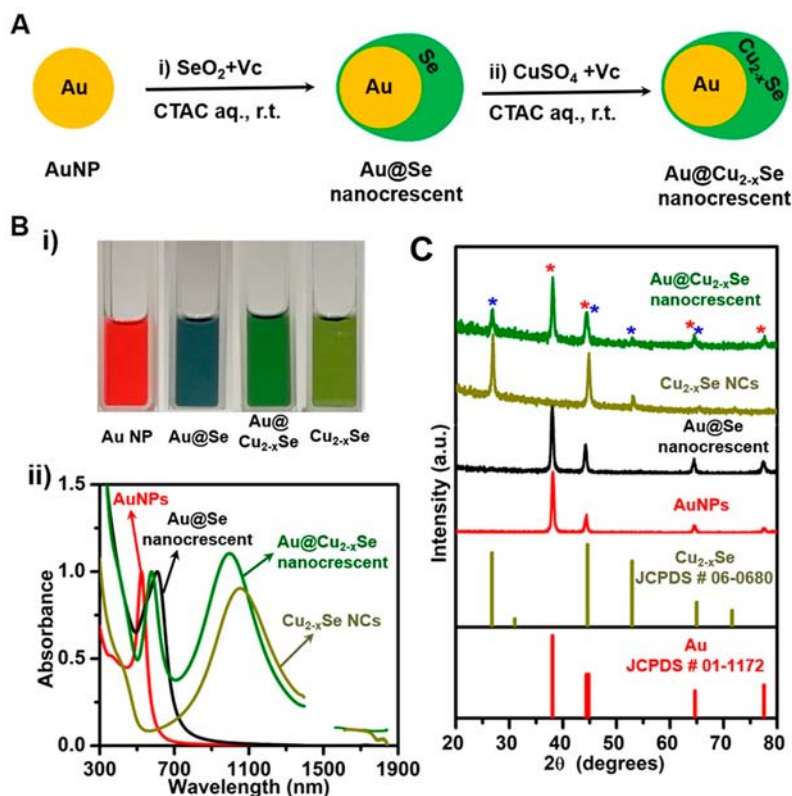


Figure 1. Synthesis and characterization of Au@Cu_{2-x}Se nanocrescents by a Se-mediated two-step process. (A) Schematic illustrating the synthesis process of Au@Cu_{2-x}Se nanocrescents by a Se-mediated two-step process: (i) the selenium layer was first coated onto the AuNPs by the reduction of SeO₂ using Vc as the reducing agent, yielding the Au@Se nanocrescents, and then (ii) the selenium layer was converted to Cu_{2-x}Se, producing the Au@Cu_{2-x}Se nanocrescents in the presence of copper(II), CTAC, and Vc. (B) (i) Optical image and (ii) optical extinction spectra of aqueous solutions of AuNPs, Cu_{2-x}Se NCs, Au@Se nanocrescents, and Au@Cu_{2-x}Se nanocrescents. The spikes arising from the strong absorption of water medium in the ~1410–1550 nm range of the extinction spectra were removed. The Au@Se nanocrescents were synthesized with 1.00 mM SeO₂, followed by the conversion into Au@Cu_{2-x}Se nanocrescents. (C) Powder XRD patterns of AuNPs, Cu_{2-x}Se NCs, Au@Se nanocrescents, and Au@Cu_{2-x}Se nanocrescents. The standard XRD patterns of Au and Cu_{2-x}Se are shown as well for reference and comparison (JCPDS no. 01-1172 of cubic Au; JCPDS no. 06-0680 of cubic berzelianite Cu_{2-x}Se).

tuned by varying the degree of the Cu vacancy. Moreover, the plasmonic properties of Cu_{2-x}Se can be postsynthetically modulated by chemically doping or electrochemically charging and discharging, a unique feature different from the classical metal plasmonics.^{20–22}

Considerable efforts have been made to develop noble metal–semiconductor hybrids that combine the respective merits of their constituents to synergistically enhance their performances in surface-enhanced spectroscopies, biosensing, bioimaging, and cancer photothermal therapy.^{23–31} The interest in such dual-plasmonic hybrids is mainly driven by the electronic interactions that modulate the energy transfer and local plasmonic field. The combination of two entirely different types of plasmonic materials such as noble metals and semiconductors could also yield new unexpected phenomena due to the interaction of two intrinsically dissimilar plasmonic building blocks. Previous studies showed that the combination of Au and Cu_{2-x}Se greatly improved the visible and NIR absorption through the strong plasmonic coupling between the Au and Cu_{2-x}Se domains, and simultaneously preserved a small size (<50 nm) suitable for the efficient cell internalization along with prolonging in vivo circulation.^{26,31–34} More importantly, the dual-plasmonic system of Au and Cu_{2-x}Se presents multiple functionalities for promising multimodality diagnostics and therapy of cancer, for example, integrating

surface-enhanced Raman spectroscopy (SERS), photoacoustic imaging, and X-ray computed tomography in a single system for diagnostics,^{23,32} along with photothermal therapy and photodynamic therapy for cancer therapy.^{35,36} However, despite intensive efforts, few reports are dedicated to elucidate the interplay between the free electron- and free hole-based plasmons in the Au and Cu_{2-x}Se domains of the Au–Cu_{2-x}Se hybrids, respectively. Previous studies showed that the plasmonic coupling caused the red-shift of the Au-associated plasmon band and the blue-shift of the Cu_{2-x}Se-associated plasmon band in the Au@Cu_{2-x}Se core–shell structures with respect to the plasmon bands of Au nanoparticles (AuNPs) and Cu_{2-x}Se nanoparticles only.^{23,32} Because of the synergistic interactions, the optical properties of the Au–Cu_{2-x}Se hybrids are not simply the sum of those of the Au and Cu_{2-x}Se constituents. The influence of the refractive index, plasmonic coupling, charge redistribution, and others could intertwine together to account for the unique plasmonic properties in such a dual-plasmonic Au–Cu_{2-x}Se system. A systematic study for the in-depth understanding of the interplay of the Au and the Cu_{2-x}Se domains in the Au–Cu_{2-x}Se hybrids becomes greatly crucial but is still lacking.

Here, we developed a selenium-mediated two-step synthetic approach to prepare the dual-plasmonic Au@Cu_{2-x}Se nanocrescents with tunable Cu_{2-x}Se domain size and NIR

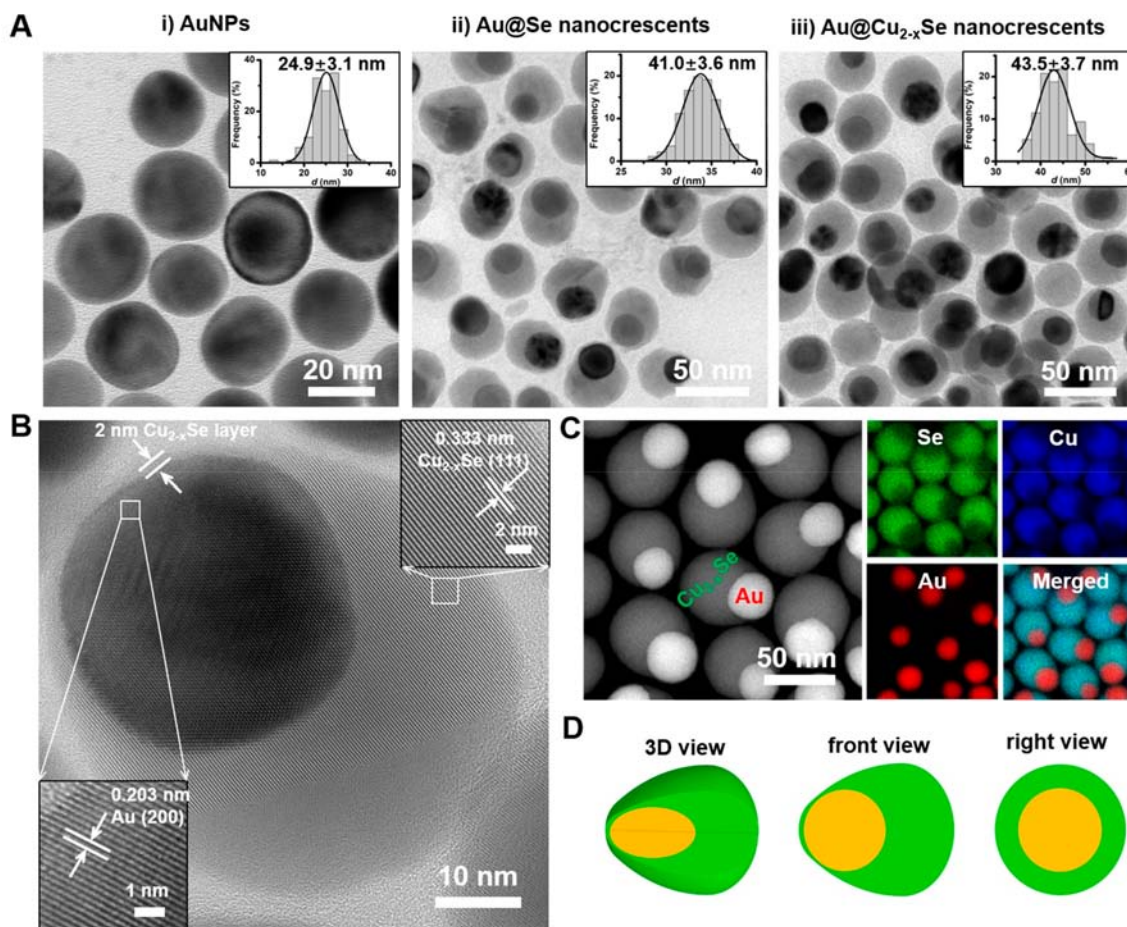


Figure 2. Morphology, size, and compositional distributions of AuNPs, Au@Se nanocrescents, and Au@Cu_{2-x}Se nanocrescents. (A) TEM images and size distribution (inset) of (i) AuNPs, (ii) Au@Se nanocrescents, and (iii) Au@Cu_{2-x}Se nanocrescents. The size of Au@Se nanocrescents and Au@Cu_{2-x}Se nanocrescents is the total size in the largest dimension of the nanoparticles containing the Au and Se or Cu_{2-x}Se domains. (B) HRTEM image of a representative Au@Cu_{2-x}Se nanocrescent. 1.0 mM SeO₂ was used for the synthesis of the Au@Cu_{2-x}Se nanocrescents. (C) HAADF-STEM and the corresponding EDX elemental mapping of the Au@Cu_{2-x}Se nanocrescent. 2.0 mM SeO₂ was used for the synthesis of the Au@Cu_{2-x}Se nanocrescents. (D) A schematic model illustrating the three-dimensional structure of the Au@Cu_{2-x}Se nanocrescent.

plasmonic properties. The Cu_{2-x}Se domain size of the nanocrescents was precisely controlled by varying the concentration of the SeO₂ precursor. The increasing Cu_{2-x}Se domain size is able to modulate the interactions between the Au and Cu_{2-x}Se domains in the Au@Cu_{2-x}Se nanocrescents and their optical properties accordingly. Unlike previous studies on the Au–Cu_{2-x}Se dimers where the interfacial contact area is partial,^{23,26,37} the complete encapsulation of the Au domain within the Cu_{2-x}Se in the present nanocrescent structures enables an intimate contact over a larger surface area without native Au surface exposure and thus excludes the unnecessary interference from the external environment. The asymmetric geometry of the nanocrescents yields the asymmetric electric field distribution along the Au–Cu_{2-x}Se interface and the nanocrescent surface, which is responsible for the unique plasmonic properties. We employed the theoretical modeling and transient absorption (TA) measurements to understand the evolution of plasmonic properties of the Au@Cu_{2-x}Se nanocrescents with the Cu_{2-x}Se domain size and elucidate the physics of the interplay between their Au and Cu_{2-x}Se domains. Results show that the electron transfer from the Au domain to the Cu_{2-x}Se domain accounts for their strong coupling interactions, modulating the carrier–phonon scattering and phonon–phonon relaxation processes. Our

findings unambiguously reveal the underlying physics of the strong coupling of the Au domain with the Cu_{2-x}Se domain in the hybrids, and provide a guidance for designing dual-plasmonic metal–semiconductor nanostructures with optimal NIR optical performances.

RESULTS AND DISCUSSION

Dual-plasmonic Au@Cu_{2-x}Se nanocrescents were synthesized by a selenium-mediated two-step process:³¹ (i) the direct synthesis of Au@Se nanocrescents and (ii) the subsequent conversion into the Au@Cu_{2-x}Se nanocrescents (Figure 1A). First, the hexadecyltrimethylammonium chloride (CTAC)-stabilized AuNPs were coated with an amorphous Se layer through the reduction of SeO₂ by the ascorbic acid (Vitamin C, Vc) reducing agent, forming the Au@Se nanocrescent structure. This synthetic method is mainly inspired by the high affinity of the Au surface toward the elemental Se and selenides, which has been widely exploited for synthesis of Au-chalcogenide core–shell NPs with large lattice mismatch.^{38,39} It should be noted that the CTAC stabilizer is critical for forming the Au@Se nanocrescent structure. The resulting Au@Se nanocrescents then were further reacted with Cu²⁺ in an aqueous solution of CTAC and Vc, yielding the Au@

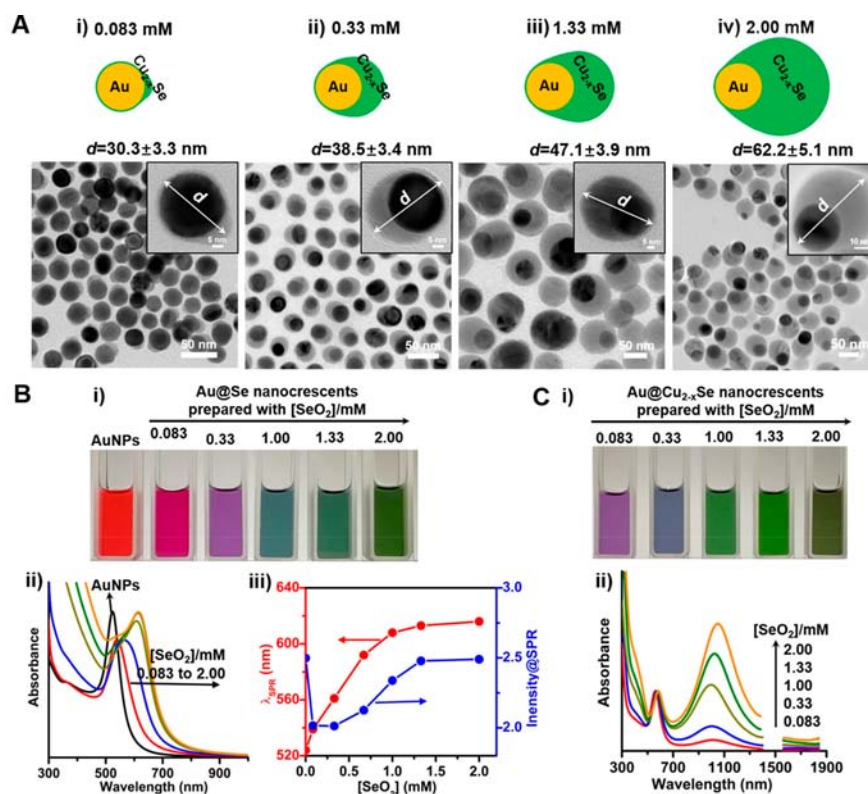


Figure 3. Evolution of the Cu_{2-x}Se domain size and optical properties with varying SeO_2 concentration of $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals. (A) TEM images (bottom) and their corresponding schematic structures (top) of $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals synthesized with the SeO_2 concentration of (i) 0.083 mM, (ii) 0.33 mM, (iii) 1.33 mM, and (iv) 2.00 mM. The insets in the panels (A(i–iv)) are the TEM images of a representative $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystal for the corresponding SeO_2 concentration, and the arrow indicates the size representing the largest dimension of the nanocrystal in their TEM images. The average size statistically obtained from the corresponding TEM images was shown as well. (B) (i) Optical image and (ii) extinction spectra of aqueous solutions of AuNPs and $\text{Au}@Se$ nanocrystals synthesized with various SeO_2 concentrations (0.083, 0.33, 1.00, 1.33, and 2.00 mM), and (iii) plot of both SPR wavelength and extinction intensity at the SPR wavelength as a function of the SeO_2 concentration. (C) (i) Optical image and (ii) extinction spectra of aqueous solutions of $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals obtained from the conversion of $\text{Au}@Se$ nanocrystals shown in (B). The spikes arising from the strong absorption of water medium in the ~ 1410 – 1550 nm range of the extinction spectra were removed.

Cu_{2-x}Se nanocrystals. Cu^{2+} was reduced to Cu^+ by the weak Vc reducing agent, and then reacted with the elemental Se to produce the Cu-deficient Cu_{2-x}Se upon the exposure to air for oxidation. The initial amorphous Se layer on the AuNP surface as the template facilitates the high control of the Cu_{2-x}Se coating onto the AuNP surface because of the large lattice mismatch between AuNPs and Cu_{2-x}Se crystals.

A wine red-to-cyan color change was accomplished immediately after SeO_2 was added to the AuNP suspension containing the Vc; the reaction solution became dark green along with the conversion of the $\text{Au}@Se$ nanocrystals into the $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals, also reflected by the spectral variation of their extinction spectra (Figure 1B(ii)). The characteristic SPR peak of the AuNPs red-shifts from 524 to 608 nm after the Se layer coating, which is ascribed to the higher refractive index ($n > 2.70$ at 524–608 nm) of Se than that ($n = 1.33$) of the water medium.^{40,41} The optical extinction spectrum of $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals made with 1.0 mM SeO_2 shows two characteristic peaks at 578 and 996 nm, which are assigned to the SPR bands corresponding to its Au and Cu_{2-x}Se domains, respectively. A blue-shift by 30 nm (from 608 to 578 nm) is observed for the Au-associated SPR band along with the conversion of the $\text{Au}@Se$ to $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals (Figure 1B(ii)). Three intertwined factors that are the refractive index change ($n = 2.6$ for Cu_{2-x}Se ; $n < 1$ for

Au at >500 nm), charge redistribution, and plasmonic coupling effects may be responsible for the red-shift of the Au-associated SPR band and the blue-shift of the Cu_{2-x}Se -associated SPR band in $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals with respect to the SPR bands of AuNPs and Cu_{2-x}Se nanocrystals (NCs), respectively. The powder XRD result confirms the amorphous form of the Se layer on the AuNP surface (Figure 1C). In the $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals, characteristic 2θ peaks at 26.7° , 44.6° , 52.9° , and 64.9° all can be indexed to the cubic berzelianite structure of Cu_{2-x}Se (JCPDS no. 06-0680 of cubic berzelianite Cu_{2-x}Se), while the 2θ peaks at 38.1° and 77.5° are attributed to the cubic Au structure (JCPDS no. 01-1172 of cubic Au). We also synthesized the Cu_{2-x}Se NCs (33.8 ± 1.9 nm) in the absence of AuNPs, showing the same cubic berzelianite structure as the $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals (Figures 1C and S3). Thus, the XRD patterns confirm the successful synthesis of highly crystalline $\text{Au}@Cu_{2-x}\text{Se}$ nanocrystals.

The TEM image clearly confirms the coating of the Se layer on the spherical AuNPs, forming the $\text{Au}@Se$ nanocrystal structure (Figure 2A). In the nanocrystal structure, the AuNP domain appearing in much darker contrast was completely encapsulated by a lighter Cu_{2-x}Se layer because of their different atomic numbers. The conversion of the amorphous Se into crystalline Cu_{2-x}Se preserves the initial nanocrystal morphology but causes a slight (~ 2 nm) increase

Table 1. Summary of Optical Properties, x Value, and Carrier Density of AuNPs, Cu_{2-x}Se NCs, and $\text{Au@Cu}_{2-x}\text{Se}$ Nanocrescents Synthesized with Various SeO_2 Concentrations

sample	$[\text{SeO}_2]$ (mM)	$N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})^a$	x value ^b	$\lambda_{\text{max}}(\text{Au})$ (nm)	$\epsilon(\text{Au})^c$ (M cm^{-1})	N_e (cm^{-3}) ^d	$\lambda_{\text{max}}(\text{Cu}_{2-x}\text{Se})$ (nm)	$\epsilon(\text{Cu}_{2-x}\text{Se})^c$ (M cm^{-1})	N_h (cm^{-3}) ^d
AuNPs	0	0	N/A	524	8.57×10^9	1.87×10^{22}	NA	NA	NA
$\text{Au@Cu}_{2-x}\text{Se}$ -1	0.083	0.18	0.20	560	2.18×10^9	1.68×10^{22}	1011	8.47×10^7	4.45×10^{21}
$\text{Au@Cu}_{2-x}\text{Se}$ -2	0.33	0.71	0.19	571	3.25×10^9	1.60×10^{22}	1002	9.25×10^8	4.37×10^{21}
$\text{Au@Cu}_{2-x}\text{Se}$ -3	1.00	1.96	0.19	578	3.00×10^9	1.58×10^{22}	996	2.37×10^9	4.30×10^{21}
$\text{Au@Cu}_{2-x}\text{Se}$ -4	1.33	2.42	0.20	579	3.46×10^9	1.54×10^{22}	1025	4.12×10^9	4.05×10^{21}
$\text{Au@Cu}_{2-x}\text{Se}$ -5	2.00	3.77	0.19	585	3.66×10^9	1.50×10^{22}	1052	6.99×10^9	3.83×10^{21}
Cu_{2-x}Se NCs	2.44	NA	0.20	NA	NA	NA	1053	1.97×10^9	4.05×10^{21}

^aThe $N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})$ ratio is the atomic ratio of the Cu_{2-x}Se to Au atoms in the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents estimated by ICP-OES. ^bThe x value was estimated by ICP-OES. ^cThe molar extinction coefficients ($\epsilon(\text{Au})$ and $\epsilon(\text{Cu}_{2-x}\text{Se})$) correspond to the SPR wavelengths $\lambda_{\text{max}}(\text{Au})$ and $\lambda_{\text{max}}(\text{Cu}_{2-x}\text{Se})$ associated with the Au and Cu_{2-x}Se domains, respectively. ^d N_e and N_h represent free electron and free hole densities estimated on the basis of the SPR bands corresponding to the Au and Cu_{2-x}Se domains, respectively.

of the total size due to the Cu atom incorporation (Figure 2A). The HRTEM image of a representative $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescent in Figure 2B unveils the complete encapsulation of the AuNP by a Cu_{2-x}Se layer of high crystallinity. The lattice fringes with the interplanar spacings of 0.333 and 0.203 nm correspond to the (111) and (200) planes of the Cu_{2-x}Se and Au domains, respectively. This is consistent with the XRD results. Both HRTEM and HAADF-STEM images clearly show that the as-prepared $\text{Au@Cu}_{2-x}\text{Se}$ hybrids possess a crescent structure where the AuNP is completely encapsulated by a Cu_{2-x}Se layer in an asymmetric manner. Furthermore, the STEM-EDX mapping shows a clear distribution of Cu and Se in the outer layer and Au in the dark core, confirming the crescent structure again. Therefore, all of this evidence corroborates the controllable synthesis of the high-quality dual-plasmonic $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents by the Se-mediated method.

We further probed the tunability of the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents by varying the SeO_2 concentration from 0.083 to 2.00 mM (Figure 3). It can be clearly seen that all samples have a crescent structure and there exists a thin (~ 2 nm) Cu_{2-x}Se layer conformally coating on the one side of the AuNP surface. The total size of the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents is proportional to the SeO_2 concentration, and the size of the Cu_{2-x}Se domain is 30.3 ± 3.3 , 38.5 ± 3.4 , 41.0 ± 3.6 , 47.1 ± 3.9 , and 62.2 ± 5.1 nm for the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents made with the SeO_2 concentrations of 0.083, 0.33, 1.00, 1.33, and 2.00 mM, respectively (Figure S4). It is worth pointing out that Cu_{2-x}Se mainly grows unidirectionally on the one side of the Au domain even with more addition of the SeO_2 precursor, forming the crescent structure. The replacement of the CTAC stabilizer with other stabilizers such as poly(vinylpyrrolidone) and poly(diallyldimethylammonium chloride) produced symmetrical core-shell or isolated nanostructures rather than the crescent-like structure. It is suggested that the shielding effect of CTAC on the deposition of Se atoms is responsible for the formation of $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents.³¹ The color of the suspension of the Au@Se nanocrescents gradually evolves from red to dark green with the SeO_2 concentration increasing up to 2.00 mM, while the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents appear violet to dark green accordingly (Figure 3B,C). The UV-vis-NIR extinction spectra show a gradual red-shift of the SPR band of the Au@Se nanocrescents with the increasing Se domain size (due to the increasing SeO_2 concentration), while the extinction intensity undergoes first a steep decrease and then a gradual

increase with the increasing Se domain size. The rationale is the relatively large refractive index of Se and the lack of the electronic interactions between the Au and the amorphous Se. There exist two SPR bands associated with the Au and Cu_{2-x}Se domains for all $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents; the SPR band corresponding to the Cu_{2-x}Se domain is located in the NIR-II region of ca. 990–1050 nm, and the SPR band corresponding to the Au domain is in the visible region of 560–585 nm, red-shifting with respect to that of the AuNPs. The intensity of the NIR SPR band is proportional to the total size of the Cu_{2-x}Se domain. We also noted the red-shift of the interband transition-induced absorption onset around 300–500 nm of Cu_{2-x}Se in $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents with the increasing Cu_{2-x}Se domain size, probably resulting from the quantum size effects (Figure 3C(ii)). Thus, we confirm that the size of the $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents can be effectively controlled by the initial concentration of the SeO_2 precursor, tuning their plasmonic properties accordingly. The corresponding plasmonic properties of all $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents along with those of AuNPs and Cu_{2-x}Se NCs are listed in Table 1.

We analyzed the elemental composition of all as-synthesized $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents by ICP-OES, and then derived their “ x ” values and the molar ratios of the Cu_{2-x}Se to Au atom accordingly ($N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})$, equivalent to the Se to Au atomic ratio) (Table 1). It can be seen that the $N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})$ value increases from 0.18 to 3.77 with the increasing SeO_2 concentration from 0.083 to 2.00 mM. Results show the similar x value of ~ 0.2 for all $\text{Au@Cu}_{2-x}\text{Se}$ nanocrescents, consistent with the literature.^{21,42} Furthermore, we determined the molar extinction coefficients ($\epsilon(\lambda)$) corresponding to the respective SPR bands of the Au and Cu_{2-x}Se domains according to the Lambert–Beer law:⁴³

$$A(\lambda) = \epsilon(\lambda) \cdot l \cdot c \quad (1)$$

where $A(\lambda)$, l , and c are the optical density at the SPR wavelength (λ), the length of the solution the incident light passes through (1 cm for the standard cuvette), and the concentration of the Au and Cu_{2-x}Se domains (mol/L), respectively. The atomic concentrations of Au, Cu, and Se determined by ICP-OES were converted to the particulate concentration of the Au and Cu_{2-x}Se domains, whose details were shown in the Supporting Information.⁴³ A dramatic decrease of the molar extinction coefficient of the Au domain-associated SPR band ($\epsilon(\text{Au})$) is observed after being encapsulated with the Cu_{2-x}Se , followed by a monotonic and

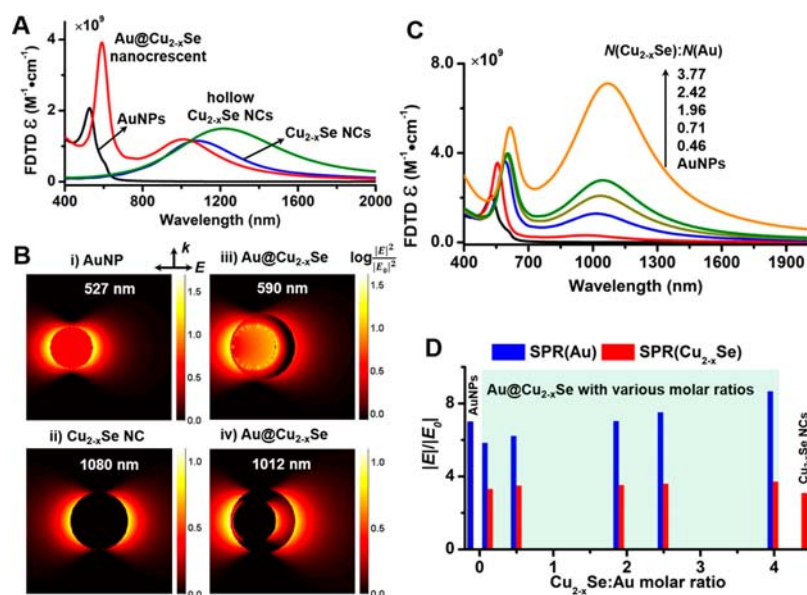


Figure 4. Theoretical analysis of plasmonic properties of Au@Cu_{2-x}Se nanocrescents with various Cu_{2-x}Se: Au molar ratios. (A) Calculated extinction spectra of AuNPs, Cu_{2-x}Se NCs, hollow Cu_{2-x}Se NCs, and Au@Cu_{2-x}Se nanocrescents. The AuNPs and Cu_{2-x}Se NCs are, respectively, modeled as a sphere of 25 and 33.8 nm in diameter taken from Figure 2, and the sizes of the AuNP and Cu_{2-x}Se domains in the nanocrescent are consistent with those of AuNPs and Cu_{2-x}Se NCs. The structure of hollow Cu_{2-x}Se NCs is similar to the Au@Cu_{2-x}Se nanocrescent except for the removal of the Au domain. (B) Distribution of electric field ($\log(|E|^2/|E_0|^2)$) of (i) AuNPs at its visible SPR wavelength (527 nm), (ii) Cu_{2-x}Se NCs at its NIR SPR wavelength (1080 nm), and (iii, iv) Au@Cu_{2-x}Se nanocrescents at its visible and NIR SPR wavelengths corresponding to (iii) the Au domain (590 nm) and (iv) the Cu_{2-x}Se domain (1012 nm). (C) Evolution of the calculated extinction spectra of Au@Cu_{2-x}Se nanocrescents with the increasing Cu_{2-x}Se: Au molar ratios. (D) The maximal electric field intensity ($|E|/|E_0|$) of AuNPs and Cu_{2-x}Se NCs at their SPR wavelengths, and of Au@Cu_{2-x}Se nanocrescents of various Cu_{2-x}Se: Au molar ratios at the SPR wavelengths corresponding to their Au and Cu_{2-x}Se domains, respectively.

slight increase with the increasing $N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})$ value; however, we can see 1–2 orders of magnitude increase of the $\epsilon(\text{Cu}_{2-x}\text{Se})$ value of the Au@Cu_{2-x}Se nanocrescents with the increasing $N(\text{Cu}_{2-x}\text{Se}):N(\text{Au})$ (Table 1).

To better understand the plasmonic properties of the dual-plasmonic Au@Cu_{2-x}Se nanocrescents, FDTD calculations of both extinction spectra and electric field distribution were performed for AuNPs, Cu_{2-x}Se NCs, and the corresponding Au@Cu_{2-x}Se nanocrescent (Figure 4A,B). The nanocrescent is modeled with an Au sphere (25 nm in diameter) asymmetrically embedded within the Cu_{2-x}Se sphere (33.8 nm in diameter), as schematically shown in Figure S2. The Au and Cu_{2-x}Se domains in the nanocrescent have the same size as the AuNPs and Cu_{2-x}Se NCs, respectively. The calculated extinction spectrum of the Au@Cu_{2-x}Se nanocrescents clearly shows two separated SPR bands located in the visible and NIR regions, associated with the Au and Cu_{2-x}Se domains, respectively (Figure 4A). This is consistent with the experimental results shown in Figure 1B. The visible SPR band of the Au@Cu_{2-x}Se nanocrescents red-shifts and concomitantly becomes stronger, while its NIR SPR band blue-shifts, with respect to those of AuNPs and Cu_{2-x}Se NCs, respectively. Surprisingly, we also observed the enhanced NIR SPR band in the hollow Cu_{2-x}Se NCs where the Au domain was removed from the Au@Cu_{2-x}Se nanocrescent. We further calculated the electric field distribution and then compared their maximal electric field intensities ($|E|/|E_0|$). The asymmetric distribution of electric field was observed in the Au@Cu_{2-x}Se nanocrescents; the maximal electric field corresponding to the visible or NIR SPR band appears at the left side of the Cu_{2-x}Se domain, quite different from those of AuNPs and Cu_{2-x}Se NCs (Figure 4A,B). A close examination

reveals that the electric field is much stronger at the right side of the Au domain of the nanocrescent than that at its left side. On the basis of the average size statistically estimated from the TEM images for all Au@Cu_{2-x}Se nanocrescents, we further examined the effect of the Cu_{2-x}Se domain size on the electric field (Figures 4C,D, S5, and S6). Both visible and NIR SPR bands gradually red-shift and concomitantly increase in the intensity with the increasing Cu_{2-x}Se domain size (Figure 4C). We can clearly see that the maximal electric field intensity at the visible SPR wavelength decreases from the AuNPs to nanocrescents, and then gradually increases with the increasing Cu_{2-x}Se domain size; in contrast, all Au@Cu_{2-x}Se nanocrescents exhibit the maximal electric field intensity at the NIR SPR wavelength much higher than Cu_{2-x}Se NCs alone, indicating the coupling effect that enhances the NIR electric field in the Au@Cu_{2-x}Se nanocrescents (Figure 4D).

According to the Drude model, we calculated the free electron (N_e) and free hole (N_h) densities corresponding to the visible and NIR SPR bands of the nanocrescents, respectively.^{18,44,45}

$$\omega_{\text{sp}} = \sqrt{\frac{\omega_p^2}{\epsilon_\infty + 2\epsilon_m} - \gamma^2} \quad (2)$$

where ω_{sp} and γ are the frequency and full-width-at-half-maximum (fwhm) of the SPR bands, respectively; ϵ_∞ and ϵ_m are the high frequency dielectric constant and the solvent dielectric constant ($\epsilon_m = n_m^2 = 1.33^2 = 1.77$ for water here), respectively; ω_p is the bulk plasma oscillation frequency related to the free electron or hole density. γ can be determined by fitting the SPR band to a Gaussian function, and ϵ_∞ can be treated as 7.0 for Cu_{2-x}Se at $x = 0.2$ according to the

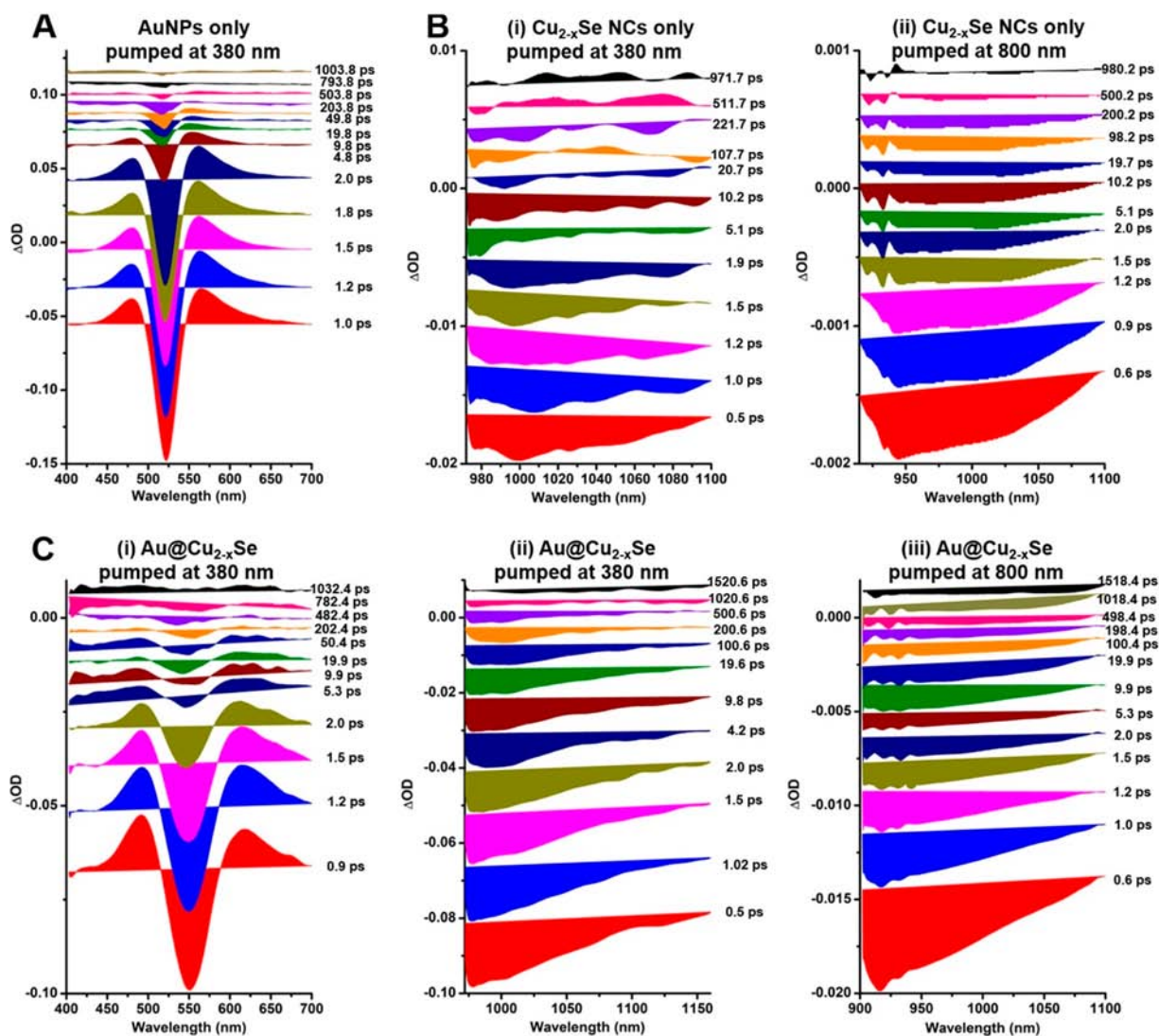


Figure 5. Transient absorption spectra of AuNPs, Cu_{2-x}Se NCs, and $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents. (A) Transient absorption spectra of AuNPs in the visible region at various delay times, pumped at 380 nm. (B) Transient absorption spectra at various delay times of Cu_{2-x}Se NCs (i) in the visible region pumped at 380 nm and (ii) in the NIR region pumped at 800 nm. (C) Transient absorption spectra at various delay times of $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents (i) in the visible region pumped at 380 nm, (ii) in the NIR region pumped at 380 nm, and (iii) in the NIR region pumped at 800 nm. The present $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents have a $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio of 0.71, which were prepared with 0.33 mM SeO_2 .

literature.⁴⁶ Thus, we calculated the free carrier density ($N = N_e$ or N_h) by the following equation:^{18,44,45}

$$\omega_p^2 = \frac{Nq_e^2}{\epsilon_0 m^*} \quad (3)$$

where m^* is the charge carrier effective mass ($m^* = m_0$ for a free electron, $m^* = 0.336m_0$ for a free hole, and m_0 is the mass of an electron),⁴⁶ q_e is the elementary charge (1.6×10^{-19} C), and ϵ_0 is the free space permittivity. From eq 3, the N_e value in the Au domain was estimated to be in the range of $(1.50\text{--}1.68) \times 10^{22} \text{ cm}^{-3}$, and the N_h value in the Cu_{2-x}Se domain is in the range of $(3.83\text{--}4.45) \times 10^{21} \text{ cm}^{-3}$; both N_e and N_h decrease with the increase of the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio from 0.18 to 3.77 (Table 1). The N_e value of all $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents is slightly lower than $1.87 \times 10^{22} \text{ cm}^{-3}$ of the AuNPs, but their N_h values are much higher than $4.05 \times 10^{21} \text{ cm}^{-3}$ of the Cu_{2-x}Se NCs except the nanocrescent of the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio of 3.77. The present $\text{Au}@Cu_{2-x}\text{Se}$

nanocrescents were made with the fixed Au domain size and the varied Cu_{2-x}Se domain size. The “ x ” value representing the degree of Cu deficiency in Cu_{2-x}Se nanoparticles is proportional to the free hole density; previous studies showed the insensitivity of the SPR wavelength of Cu_{2-x}Se nanoparticles to the particulate size. As shown in Table 1, all $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents possess the similar “ x ” value of 0.2, which cannot explain the different free hole densities among these nanocrescents. We attribute the different hole densities to the interplay between the Au domain and the Cu_{2-x}Se domain due to the varied Cu_{2-x}Se domain size.

TA spectra have been widely used to study the carrier dynamics, and the carrier–phonon and phonon–phonon interactions in plasmonic noble metals, but few efforts have been devoted to examine the dynamics of charge carriers in the $\text{Au}-\text{Cu}_{2-x}\text{Se}$ hybrid system. Here, we performed TA studies to gain insights into the carrier dynamics in the $\text{Au}@Cu_{2-x}\text{Se}$ nanocrescents pumped with a 100 fs pulse laser at 380 and 800 nm, respectively (Figure 5). Because the threshold energy of

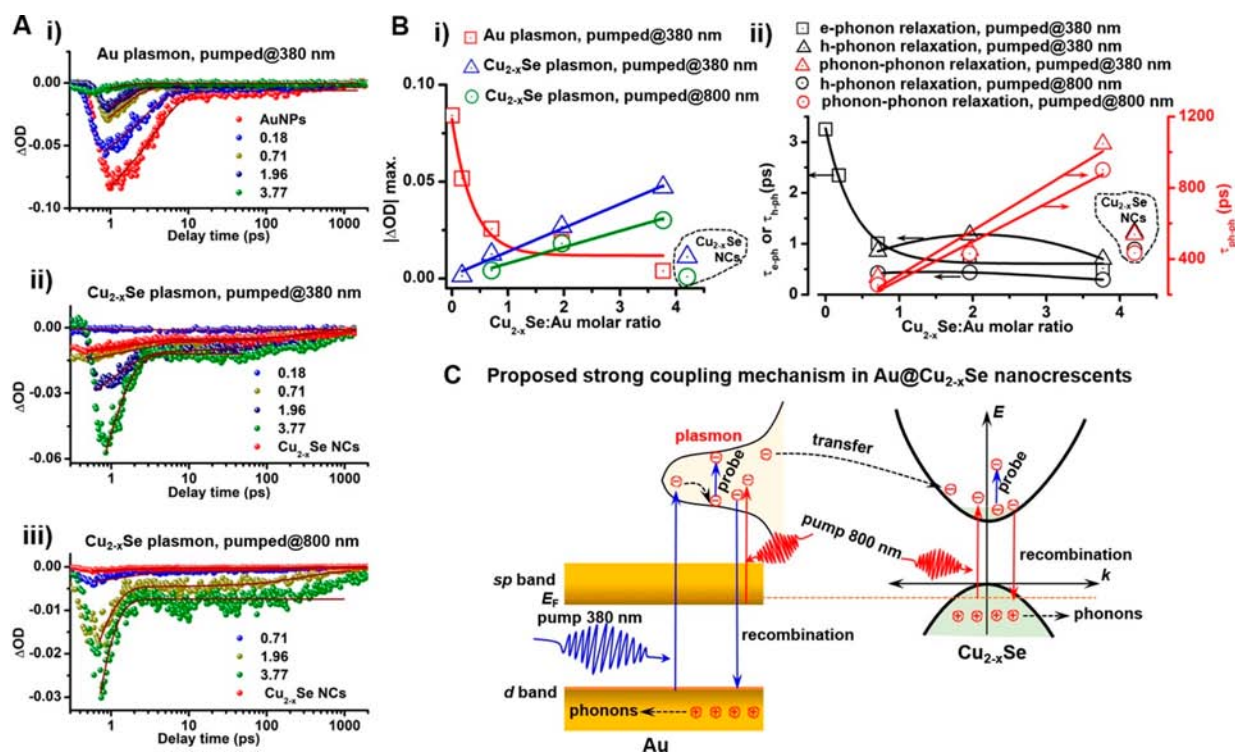


Figure 6. Cu_{2-x}Se:Au molar ratio-dependent bleaching dynamics and proposed coupling mechanism in Au@Cu_{2-x}Se nanocrystals. (A) Bleaching dynamics of the TA signal (i) of the intrinsic SPR of AuNPs and the Au domain in Au@Cu_{2-x}Se nanocrystals of various Cu_{2-x}Se:Au molar ratios pumped at 380 nm, and (ii, iii) of the intrinsic NIR SPR of Cu_{2-x}Se NCs and the Cu_{2-x}Se domain in Au@Cu_{2-x}Se nanocrystals of various Cu_{2-x}Se:Au molar ratios pumped at (ii) 380 nm and (iii) 800 nm. (B) Effects of Cu_{2-x}Se:Au molar ratio on (i) the bleaching amplitude and (ii) electron/hole/photon–phonon relaxation time constant at the SPR wavelengths of the Au and Cu_{2-x}Se domains pumped at 380 and 800 nm, respectively. Results of AuNPs and Cu_{2-x}Se NCs are shown as well. (C) Schematic illustrating the proposed coupling mechanism in the dual-plasmonic Au@Cu_{2-x}Se nanocrystals.

the interband transition in AuNPs is ~ 2.4 eV (~ 520 nm), the pulse laser of 380 nm is able to trigger its interband transition from 5d-band to the hybridized 6sp-band, resulting in a transient increase in both hole density in the 5d-band and electron density in the 6sp-band.⁴⁷ With the excitation of 380 nm-pulse laser, the TA spectra of AuNPs exhibit a negative bleaching band at ~ 528 nm and two positive absorption bands centered at 480 and 565 nm, respectively (Figure 5A). The negative TA signal is due to the “bleaching” of the intrinsic Au SPR band; the positive TA signals at 480 and 565 nm are due to the thermal broadening of the SPR band resulting from an increase in the electronic temperature induced by the pulse laser excitation.⁴⁸ The TA spectra of Cu_{2-x}Se NCs show a broad negative band located in the NIR (900–1100 nm) region originating from the bleaching of the intrinsic NIR SPR of Cu_{2-x}Se NCs, observable under the 380 nm- or 800 nm-pulse laser. Previous studies suggested that, with the excitation of the pulse laser, Cu_{2-x}Se NCs followed bleaching behaviors of plasmonic electrons similar to those in noble metal (Au and Ag) nanoparticles.²¹ Time-resolved TA spectra show the gradual decrease of the TA signals of the intrinsic SPRs of AuNPs and Cu_{2-x}Se NCs for both 380 nm- and 800 nm-pulse lasers (Figure 5A,B). In addition, we found that the intrinsic TA signal of Cu_{2-x}Se NCs under the 800 nm-pulse laser was much lower than that under the 380 nm-pulse laser, indicating the additional contribution from the resonant excitation with the 380 nm-pulse laser. Similarly, both positive and negative TA bands of the Au domain of the Au@Cu_{2-x}Se nanocrystal with the Cu_{2-x}Se:Au molar ratio of 0.71 weaken over delay

time as well, but their TA signals are much weaker than those of AuNPs (Figure 5C). In addition, we surprisingly found that the intrinsic TA signal of its Cu_{2-x}Se domain was much stronger in the Au@Cu_{2-x}Se nanocrystals than that in Cu_{2-x}Se NCs under both 380 and 800 nm pulse lasers. This indicates the enhanced NIR optical absorption due to the electronic interplay between the plasmonic Au and Cu_{2-x}Se domains of the Au@Cu_{2-x}Se nanocrystal.

We further examined the bleaching dynamics of the intrinsic TA signals of Au@Cu_{2-x}Se nanocrystals of different Cu_{2-x}Se:Au molar ratios (Figure 6A,B). The excitation with the pulse laser creates a nonthermal distribution of charge carriers (electrons in the Au domain and holes in the Cu_{2-x}Se domain) and accordingly causes an increase in the carrier temperature, producing a thermalized electron gas (hot electrons) with higher temperature than the lattices.^{21,49,50} The newly generated distribution of carriers can be reached on a time scale of subpicosecond via strong carrier–carrier scattering, followed by the bleaching of the SPR bands through fast carrier–phonon coupling within few picoseconds and slow phonon–phonon coupling within a few hundred picoseconds. It can be seen that the bleaching amplitude of the intrinsic TA signal of the Au domain exponentially decreases with the increase of the Cu_{2-x}Se:Au molar ratio under the excitation of 380 nm-pulse laser, while the bleaching amplitude of the intrinsic TA signal of the Cu_{2-x}Se domain linearly increases with the increasing Cu_{2-x}Se:Au molar ratio under both 380 nm- and 800 nm-pulse lasers, except the Au@Cu_{2-x}Se nanocrystal with the Cu_{2-x}Se:Au molar ratio of 0.18

due to the low loading (accordingly too weak NIR SPR absorption) of Cu_{2-x}Se (Figure 6B(i)). It is worth noting that the bleaching amplitude of the TA signal of the Cu_{2-x}Se SPR in the nanocrescents is much larger than the Cu_{2-x}Se NCs over the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio range of 0.71–3.77 regardless of the excitation of the 380 nm- or 800 nm-pulse laser. These results suggest the decreasing electron density in the Au domain due to the transfer of electrons from the Au domain to the Cu_{2-x}Se domain; the increasing size of the Cu_{2-x}Se domain could increase the NIR absorption cross-section and thus compensate the neutralization effect of electron transfer to the hole-rich Cu_{2-x}Se domain, leading to an increase in the bleaching amplitude with the increasing $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio.

Furthermore, we found that the dynamic TA signal of the intrinsic SPR of Au followed a monoexponential decay over time, while the dynamic TA signal of the intrinsic SPR of Cu_{2-x}Se showed a biexponential decay process (Figure 6A). Thus, the fast carrier–phonon scattering time constant ($\tau_{\text{c-ph}}$) and the slow phonon–phonon relaxation time constant ($\tau_{\text{ph-ph}}$) of $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents can be extracted from the function as follows:^{48,51}

$$f(t) = A_f \exp\left(-\frac{t}{\tau_{\text{c-ph}}}\right) + A_s \left(-\frac{t}{\tau_{\text{ph-ph}}}\right) + A_0 \quad (4)$$

where the first term corresponds to the carrier–phonon scattering process, and the second term and the constant term (A_0) represent the phonon–phonon relaxation process; A_f and A_s are the corresponding constant coefficients for the first and second terms, respectively. The fitting of the dynamic TA curve was performed in the range from the delay time with the SPR bleaching maximum to the delay time when the TA signal is almost constant over time. The electron–phonon scattering time constant ($\tau_{\text{e-ph}}$) and hole–phonon scattering time constant ($\tau_{\text{h-ph}}$) can be obtained from the dynamic curves of the TA signals of the SPRs of Au and Cu_{2-x}Se , respectively. Because Cu_{2-x}Se has a carrier behavior similar to that of noble metals,^{21,49} we assume that the relaxation processes of excited carriers in Cu_{2-x}Se have the same physical origin as those of noble metals. Thus, we can achieve two time-constants from the fitting results and assign the small and large time-constants as the $\tau_{\text{c-ph}}$ and $\tau_{\text{ph-ph}}$, respectively. The monoexponential decay of the TA signal of the Au SPR indicates the constant phonon–phonon relaxation time constant involved in the constant term (A_0).⁵¹ Results show that the $\tau_{\text{e-ph}}$ value exponentially decreases with the increasing $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio, and becomes almost constant when the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio exceeds 0.71; the $\tau_{\text{h-ph}}$ value that varies insignificantly with the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio is much larger under the 380 nm-pulse laser than that under the 800 nm-pulse laser over the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ ratio range of 0.71–3.77. In addition, the hole–phonon relaxation of the Cu_{2-x}Se SPR in the nanocrescents is much faster than Cu_{2-x}Se NCs over the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio range of 0.71–3.77 under both 380 nm- and 800 nm-pulse lasers. The $\tau_{\text{ph-ph}}$ value linearly increases with the increase of the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio over the range investigated for both 380 nm- and 800 nm-pulse lasers. The $\tau_{\text{ph-ph}}$ value of Cu_{2-x}Se NCs is much larger than those of the nanocrescents with the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio of 0.71 and 1.96, but smaller than that of the nanocrescents with the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio of 3.77, regardless of the 380 nm- or 800 nm-pulse laser. Unfortunately, we were unable to obtain

the $\tau_{\text{h-ph}}$ and $\tau_{\text{ph-ph}}$ values for $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents of a $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio of less than 0.71 because of the low TA signal of the intrinsic SPR of Cu_{2-x}Se . It is worth noting that both bleaching amplitude and $\tau_{\text{e-ph}}$ of the Au SPR follow the similar change trend with the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio. The rationale is that the electron transfer from the Au domain to the Cu_{2-x}Se domain decreases electron density and thus accelerates the electron–phonon relaxation in the Au domain. Similarly, both the bleaching amplitude of the Cu_{2-x}Se SPR and the phonon–phonon relaxation time-constant increase with the $\text{Cu}_{2-x}\text{Se}:\text{Au}$ molar ratio. The electron transfer from the Au domain to the Cu_{2-x}Se domain decreases the hole density in the Cu_{2-x}Se domain, but the increasing Cu_{2-x}Se domain size results in the increase of the extinction cross-section of the NIR SPR band, which slows the phonon–phonon relaxation process.

On the basis of the aforementioned analysis, we propose a possible mechanism that elucidates the interactions between the Au and Cu_{2-x}Se domains of the $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents, schematically illustrated in Figure 6C. Both interband transition and SPR excitation would induce the generation of hot electrons, which could subsequently transfer from the Au domain to the Cu_{2-x}Se domain. The increasing Cu_{2-x}Se domain size enhances the extraction of hot electrons from the Au domain, which decreases the electron density in the Au domain. Despite the charge neutralization in the Cu_{2-x}Se domain, the increasing Cu_{2-x}Se size causes the increase of the NIR extinction cross-section and thereby increases the bleaching amplitude of the TA signal of the intrinsic SPR of Cu_{2-x}Se . The hot electrons and holes transfer to the surrounding lattices through the electron–phonon and hole–phonon interactions, which enhance the phonon–phonon relaxation process, and weaken both the electron–phonon scattering in the Au domain and the hole–phonon scattering in the Cu_{2-x}Se domain of the nanocrescents. Therefore, this work presents new physical insights into the coupling interactions existing in the $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents, and thereby an in-depth understanding of their optical properties, benefiting the design of Au– Cu_{2-x}Se hybrids of optimal optical properties for a diversity of plasmon-based applications.

CONCLUSIONS

In summary, a selenium-mediated two-step synthetic approach has been developed to prepare high-quality dual-plasmonic $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents, which allows tuning the Cu_{2-x}Se domain size and the plasmonic properties in the NIR-II region by changing the concentration of the SeO_2 precursor. Cu_{2-x}Se grows unidirectionally on the Au core with more addition of the SeO_2 precursor, forming the asymmetric crescent structure responsible for the unique plasmonic properties of $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents. The $\text{Au}@\text{Cu}_{2-x}\text{Se}$ nanocrescents exhibit two characteristic SPR bands positioned at the visible and NIR-II regions that originate from the collective oscillation of free electrons and free holes in the Au and Cu_{2-x}Se domains of the nanocrescents, respectively. With the increasing Cu_{2-x}Se domain size, the NIR SPR band of the Cu_{2-x}Se domain becomes more intense with a slight red-shift. The theoretical modeling shows the asymmetric electric field distributions arising from the asymmetric structure of the nanocrescents. Both experimental results and theoretical calculation confirm the strong coupling between the Au and Cu_{2-x}Se domains of the nanocrescents, accounting for their enhanced optical

properties with respect to AuNPs and Cu_{2-x}Se NCs. TA measurements suggest the Cu_{2-x}Se domain size-dependent bleaching amplitude and carrier dynamics of the intrinsic SPRs corresponding to the Au and Cu_{2-x}Se domains with variations of the carrier–phonon and phonon–phonon coupling interactions. A coupling mechanism between the Au and Cu_{2-x}Se domains of the nanocrescents is proposed to explain the Cu_{2-x}Se-dependent plasmonic properties. This work presents not only a synthetic method for the high-quality dual-plasmonic Au@Cu_{2-x}Se nanocrescents but also generates new physical insights into the interplay between their Au and Cu_{2-x}Se domains. We foresee a huge potential of this type of dual-plasmonic hybrids for NIR-II plasmonic applications in photovoltaics, photocatalysis, and biomedicine.

MATERIALS AND METHODS

Chemicals and Materials. All chemicals and solvents were of analytical grade and used as received unless otherwise specified. Sodium borohydride (NaBH₄, >98%), copper(II) sulfate pentahydrate (CuSO₄·5H₂O, ≥99%), ascorbic acid (Vc, ≥99%), and selenium dioxide (SeO₂, ≥99%) were obtained from Sinopharm Chemical Reagent Co., Ltd. Chloroauric acid (HAuCl₄·4H₂O, 99% trace metals basis) was ordered from Shanghai Civi Chemical Technology Co., Ltd. Hexadecyltrimethylammonium chloride (CTAC, 97%) was purchased from Aladdin Chemistry Co., Ltd. Gold standard solution (1000 μg/mL) containing 1.5 M HCl, copper standard solution (1000 μg/mL) containing 1.0 M HNO₃, and selenium standard solution (1000 μg/mL) containing 2.0 M HNO₃ were obtained from Guobiao (Beijing) Testing and Certification Co., Ltd. Ultrapure water with a resistivity of 18.2 MΩ cm (at 25 °C) was produced with a Millipore Direct-Q3 UV system (Molshein, France) and used in the entire experiment. All glassware used in this work was cleaned with aqua regia, rinsed thoroughly with ultrapure water, and then air-dried prior to use throughout all experiments.

Synthesis of Au Nanoparticles (AuNPs). The spherical AuNPs were synthesized according to the well-established protocol described in the literature.⁵² First, 0.3 mL of 10 mM ice-cold NaBH₄ aqueous solution was added to a solution containing 2.5 mL of 0.15 M CTAC solution and 0.025 mL of 50 mM HAuCl₄. After being vigorously stirred for 2 min, the solution was kept still at room temperature for 1 h, producing the Au seed solution. Next, 3.2 mL of 0.2 M CTAC solution was mixed with 0.8 mL of 10 mM HAuCl₄ solution at room temperature, followed by successive addition of 35 mL of ultrapure water and 3.8 mL of 0.1 M Vc solution under gentle stirring, producing the growth solution. Ten microliters of the Au seed solution was added to the growth solution at room temperature. After being vigorously shaken for 1 min, the reaction solution was aged at 30 °C for at least 8 h. Finally, the reaction solution was centrifuged at 7500 rpm, and the pellet was the AuNP products with an optical extinction peak of 524 nm, which were dispersed in ultrapure water with a final concentration of 2.0 nM for further use.

Synthesis of Cu_{2-x}Se Nanocrystals. 200 μL of 0.1 M Vc aqueous solution was added into 3 mL of 20 mM CTAC aqueous solution followed by the addition of SeO₂ with a final concentration of 2.44 mM. After being stirred for 10 min, a freshly prepared mixture consisting of 80 μL of 0.2 M CuSO₄·5H₂O and 0.4 mL of 0.1 M Vc aqueous solution was added under vigorous stirring. The reaction was allowed to continue until a dark green solution was observed, indicating the formation of Cu_{2-x}Se NCs.

Synthesis of Au@Cu_{2-x}Se Nanocrescents. Two milliliters of the 20 mM CTAC aqueous solution was added to 1 mL of 2.0 nM AuNP suspension under vigorous stirring at 30 °C. Next, 200 μL of 0.1 M Vc aqueous solution was added, followed by the addition of a given amount of SeO₂ aqueous solution to get a final SeO₂ concentration of 0.083, 0.33, 1.00, 1.33, or 2.00 mM. After being stirred for 10 min, a solution containing 30 μL of 0.2 M CuSO₄·5H₂O and 0.4 mL of 0.1 M Vc aqueous solution was added. The mixture was allowed to react under vigorous stirring at room temperature for

overnight. The resulting products were purified by successive centrifugation and washing, and then redispersed in 1 mL of ultrapure water for further use.

Characterization and Instrumentation. Transmission electron microscopy (TEM) images were taken using an FEI Tecnai G2 F20 transmission electron microscope (FEI, U.S.) operating at 200 kV; high-resolution TEM (HRTEM) images, high-angle annular dark-field scanning TEM (HAADF-STEM) images, and energy dispersive X-ray (EDX) elemental mapping were obtained on an FEI Titan G2 80-200 transmission electron microscope equipped with a spherical aberration corrector, allowing it to reach a spatial resolution of ~0.7 Å at 300 kV. The geometric parameters of all samples were measured from the TEM data using the ImageJ analysis software. At least 100 nanoparticles were counted for each type of sample, and Gaussian fitting was carried out to obtain the average size of nanoparticles. UV–vis–NIR extinction spectra were measured in the wavelength range from 175 to 3300 nm using a Cary 5000 UV–vis–NIR spectrophotometer (Agilent Technology, U.S.). The crystal structures were determined by a SmartLab3 kW X-ray diffractometer (Rigaku, Japan) using Cu Kα radiation (λ = 1.5406 Å) operating at 40 kV and 30 mA. The elemental composition of samples was measured by inductively coupled plasma–optical emission spectroscopy (ICP–OES) on an Agilent 5100 inductively coupled plasma–optical emission spectrometer (Agilent Technologies, U.S.). Samples for the ICP–OES measurements were digested in concentrated nitric acid, followed by dilution with ultrapure water.

Transient Absorption Measurements. The TA measurements were performed at room temperature on the ExciPro femtosecond pump–probe system (CDP Systems Corp.), equipped with a regeneratively amplified Ti:sapphire ultrafast laser system for the excitation (Coherent Astrella).^{53,54} The incident pump pulses with the wavelength of 380/800 nm and pulse width of 100 fs were generated by an ultrafast optical parametric amplifier (Coherent OperA Solo). The probe beams were produced from a thin CaF₂ plate, which were focused into a fiber-coupled multichannel spectrometer with complementary metal oxide semiconductor sensors at a frequency of 1 kHz. The power of the pump pulses was adjustable with a series of neutral density filter wheels. The delay between the pump and probe pulses was controlled by a motorized delay stage. The TA spectra and dynamic data were collected over a broad spectral range in a transmission mode with a 0.3 mW pump power and a ~300 μm focused spot size of pumping beam on the aqueous suspension of samples, which were placed in a 10 mm-thick quartz cuvette. The concentration of the samples was controlled to have the same power density for the excitation, which allows for the comparison of the bleaching amplitude between samples.

Theoretical Modeling of the Plasmonic Properties. Both extinction spectra and electric field distribution were calculated by the Finite-Difference Time-Domain (FDTD) method using the commercial software Lumerical FDTD Solutions 8.19.1584 (Lumerical Solutions, Inc., Canada). For the simplicity of Au@Cu_{2-x}Se nanocrescents, we modeled an AuNP core with a sphere of 25 nm in diameter, conformally coated with a 2 nm-thick Cu_{2-x}Se layer and then deposited with an additional asymmetric Cu_{2-x}Se shell of various sizes, schematically shown in Figure S2. The dimensional parameters of the Au and Cu_{2-x}Se domains were statistically measured from the TEM data and used for the model creation. The dielectric function of Au was taken from the measured data by Johnson and Christy,⁵⁵ while the dielectric function (ε(ω) = ε₁(ω) + iε₂(ω)) of Cu_{2-x}Se with various Cu_{2-x}Se domain sizes was calculated by the Drude model:^{18,44,45}

$$\epsilon_1 = \epsilon_\infty - \frac{\omega_p^2}{\omega^2 + \gamma^2} \quad (5)$$

$$\epsilon_2 = \frac{\omega_p^2 \gamma}{\omega(\omega^2 + \gamma^2)} \quad (6)$$

where ε_∞ is the high frequency dielectric permittivity, and ω_p and γ are the plasma frequency of free carriers of the nanoparticles and the

free carrier damping, respectively. The background dielectric constant ϵ_∞ is due to the presence of interband transitions at higher photon energy, for example, $\epsilon_\infty = 1$ for Au, $\epsilon_\infty = 7$ for Cu_{2-x}Se at $x \approx 0.2$.⁴⁶ The background refractive index of the surrounding medium (water) was set to be 1.33. A total-field scattered-field was employed as the incident source for calculations of both electric field distribution and extinction spectra in the wavelength range from 400 to 2000 nm. The mesh size of the box containing a geometric structure was set as 0.25 nm. The calculated extinction cross-section (σ_{theory}) from the FDTD simulations was converted to the theoretical molar extinction coefficients (ϵ_{theory}) by the following equation:⁵⁶

$$\epsilon_{\text{theory}} = \frac{N_A}{2.3 \times 10^3} \cdot \sigma_{\text{theory}} \quad (\text{M}^{-1} \text{cm}^{-1}) \quad (7)$$

(where N_A is Avogadro's number.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.chemmater.9b04100>.

Calculation of the concentration of the particulate size of nanoparticles, unit cell structures of cubic Au and cubic berzelianite Cu_{2-x}Se , crystallographic parameters of the unit cells of cubic Au and cubic berzelianite Cu_{2-x}Se , schematic of the model of the $\text{Au}@ \text{Cu}_{2-x}\text{Se}$ nanocrescents for the FDTD simulations, TEM image and size distribution of as-synthesized Cu_{2-x}Se NCs, size distribution of $\text{Au}@ \text{Cu}_{2-x}\text{Se}$ nanocrescents, and the Cu_{2-x}Se domain size-dependent electric field ($\log|E|^2/|E_0|^2$) distributions at the calculated SPR wavelengths corresponding to the Au and Cu_{2-x}Se domains (PDF)

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Notes

The authors declare no competing financial interest.

■ ACKNOWLEDGMENTS

We would like to acknowledge financial support by the National Natural Science Foundation of China (Grants 51871246 and 11574130), the Innovation-Driven Project of Central South University (Grant 2018CX002) and Hunan Provincial Science & Technology Program (Grant 2017XK2027), and the Shenzhen Science and Technology Innovation Commission (Grants KQJSCX20170726145748464 and JCYJ20180305180553701).

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